



Holomorphic functions and subelliptic heat kernels over lie groups.

Later parts of the talk describes work with Leonard Gross, Laurent Saloff-Coste.

Bruce Driver

Department of Mathematics, 0112
University of California at San Diego, USA
<http://math.ucsd.edu/~driver>

GEOQUANT

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Some papers

- D., Gross, Saloff-Coste. Holomorphic functions and subelliptic heat kernels over Lie groups. *J. European Mathematical Society*, **11**, 941—978 (2009).
- D., Gross, Saloff-Coste. Surjectivity of the Taylor map for complex nilpotent Lie groups. *Math. Proc. Camb. Phil. Soc.*, **146**, 177–195. (2009)
- D., Gross, Saloff-Coste. Growth of Taylor coefficients over complex homogeneous spaces. *Preprint*.

Basic Notation

- (G, o) = a pointed complex manifold
- $\mathcal{H}(G) := \{f : G \rightarrow \mathbb{C} \mid f \text{ is holomorphic}\}$
- λ be a measure on G ,

$$(f, g)_{L^2(\lambda)} := \int_G f \cdot \bar{g} \, d\lambda$$

$$\|f\|_{L^2(\lambda)}^2 := (f, f)_{L^2(\lambda)} = \int_G |f|^2 \, d\lambda$$

$$\mathcal{H}L^2(G) := \left\{ f \in \mathcal{H}(G) : \|f\|_{L^2(\lambda)} < \infty \right\}$$

Fock, Kakutani, Itô, Segal, Bargmann, Gross, Hall

- Suppose G = Compact type Lie group and $o = e$
- \mathcal{D}_t is a space of derivatives of holomorphic functions
- $Tf := \{\text{“derivatives” of } f \text{ at } e\}$

$$\begin{array}{ccc}
 L^2(G, p_t) & \xrightarrow{T \circ e^{t\Delta/2}} & \mathcal{D}_t \\
 \searrow & & \nearrow \\
 \mathcal{A} \circ e^{t\Delta/2} & \mathcal{HL}^2(G^{\mathbb{C}}, \mu_t) & T
 \end{array} \tag{1}$$

Two Basic Questions

Let

$$\mathcal{D} := \{\text{“derivatives” of } f \text{ at } o : f \in \mathcal{H}(G)\}$$

be the **derivative space** associated to $\mathcal{H}(G)$ and let $T : \mathcal{H}(G) \rightarrow \mathcal{D}$ be the **“Taylor map;”**

$$Tf := \{\text{“derivatives” of } f \text{ at } o\}.$$

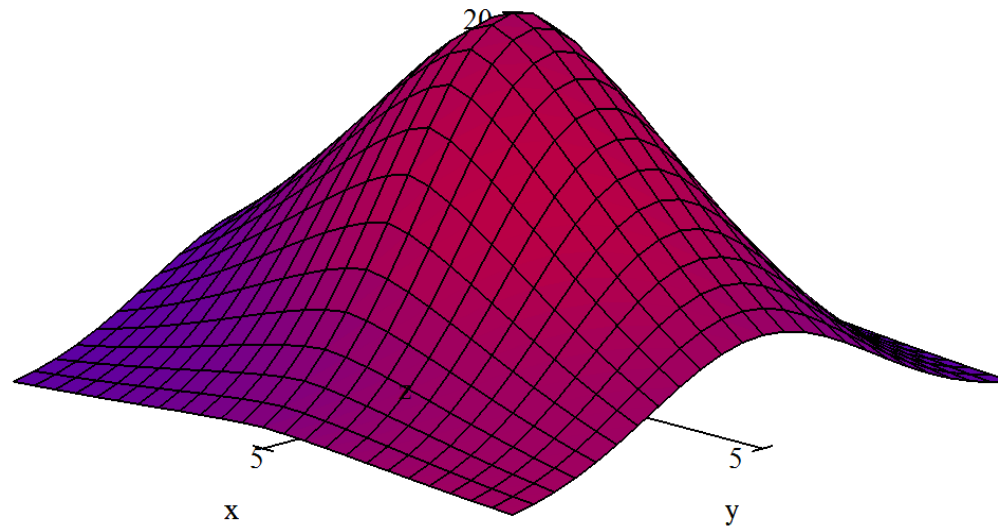
1. Characterize the derivative space, \mathcal{D} .
2. Find the norm, $\|\cdot\|_{\mathcal{D}}$, on \mathcal{D} such that

$$\int_G |f|^2 d\lambda = \|Tf\|_{\mathcal{D}}^2 \text{ for all } f \in \mathcal{H}(G).$$

- We will begin with the case, $G = \mathbb{C}$ and $o = 0$.

The Case $G = \mathbb{C}$

- Let $G = \mathbb{C}$ and $o = 0$
- $z = x + iy$
- $dm(z) = dx dy$
- $d\lambda = \rho(z) dm(z)$ with $\rho \in C(\mathbb{C}, (0, \infty))$.



$$\rho \in C(\mathbb{C}, (0, \infty))$$

Notation (Taylor map). Given a function, f , which is holomorphic near 0, let

$$Tf = \left\{ f^{(n)}(0) \right\}_{n=0}^{\infty} \in \mathbb{C}^{\mathbb{N}_0}.$$

Notation (Derivative Space). Let

$$\mathcal{D} := \left\{ \alpha := \{\alpha_n\}_{n=0}^{\infty} \subset \mathbb{C} : \limsup_{n \rightarrow \infty} \left| \frac{\alpha_n}{n!} \right|^{1/n} = 0 \right\}.$$

Theorem 1 (Taylor's Theorem & Root Test). *If $f \in \mathcal{H}(\mathbb{C})$ then*

- **The Root Test:** $Tf \in \mathcal{D}$,
- **Taylor's Theorem:** $T : \mathcal{H}(\mathbb{C}) \rightarrow \mathcal{D}$ is a linear isomorphism with inverse,

$$T^{-1}(\alpha)(z) = \sum_{n=0}^{\infty} \frac{\alpha_n}{n!} z^n \text{ for all } z \in \mathbb{C}.$$

Goals

1. Develop some basic properties of $\mathcal{H}L^2(\lambda)$.
2. Identify the norm on \mathcal{D} which makes

$$T|_{\mathcal{H}L^2(\lambda)} : \mathcal{H}L^2(\lambda) \rightarrow \mathcal{D} \quad \text{isometric.}$$

3. Characterize the image, $T(\mathcal{H}L^2(\lambda)) \subset \mathcal{D}$, of the Taylor map.

Complex Analysis Basics

- $\partial_x := \frac{\partial}{\partial x}$ and $\partial_y := \frac{\partial}{\partial y}$.
- $\partial := \partial_z = \frac{1}{2} \left(\frac{\partial}{\partial x} - i \frac{\partial}{\partial y} \right)$ and $\bar{\partial} := \partial_{\bar{z}} = \frac{1}{2} \left(\frac{\partial}{\partial x} + i \frac{\partial}{\partial y} \right)$ so
- $\partial_x = \partial + \bar{\partial}$ and $\partial_y = i(\partial - \bar{\partial})$.
- If $\Delta := \partial_x^2 + \partial_y^2$, then

$$\Delta = (\partial + \bar{\partial})^2 - (\partial - \bar{\partial})^2 = 4\partial\bar{\partial}.$$

Corollary 2. *If $f \in \mathcal{H}(\mathbb{C})$, then $|f|^2$ is sub-harmonic,*

$$\Delta |f|^2 = 4|\partial f|^2 = 4|\partial_x f|^2 \geq 0.$$

Proof: The Cauchy Riemann equations imply,

$$\bar{\partial}f = 0 \text{ and } \partial\bar{f} = 0$$

and therefore,

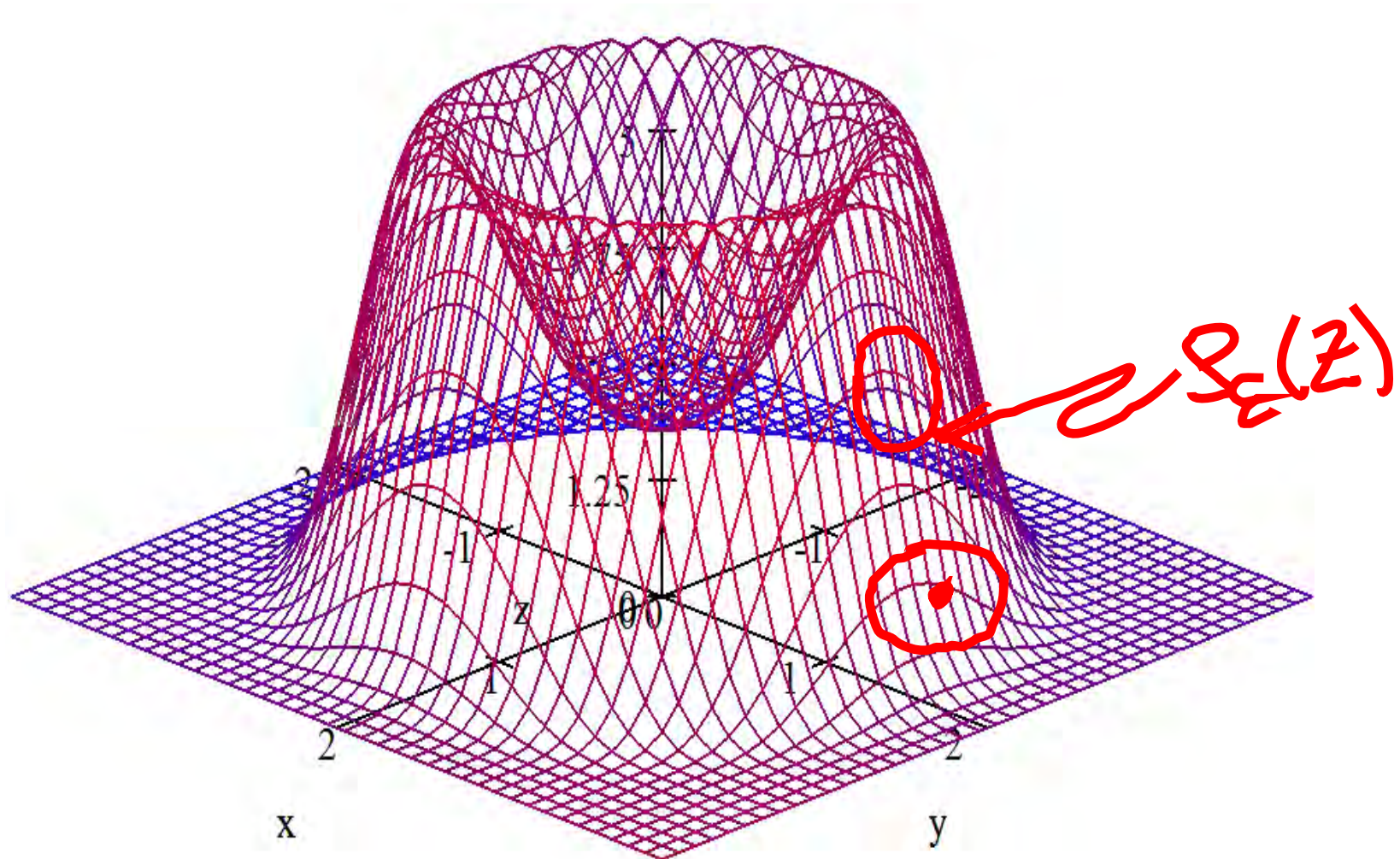
$$\Delta |f|^2 = 4\partial\bar{\partial}(f\bar{f}) = 4\partial f\bar{\partial}\bar{f}.$$

Q.E.D.

Pointwise Bounds

Notation. For every $\varepsilon > 0$, let

$$\rho_\varepsilon(z) := \min \{ \rho(w) : w \in D(z, \varepsilon) \}.$$



Theorem 3 (Crude Pointwise Bounds). *Suppose that $g \geq 0$ is a sub-harmonic (i.e. $\Delta g \geq 0$), then*

$$g(z) \leq \|g\|_{L^1(\lambda)} \frac{1}{\pi \varepsilon^2} \frac{1}{\rho_\varepsilon(z)} \quad \forall z \in \mathbb{C}. \quad (2)$$

In particular if $f \in \mathcal{H}(\mathbb{C})$, then

$$|f(z)|^2 \leq \frac{1}{\pi \varepsilon^2} \|f\|_{L^2(\lambda)}^2 \frac{1}{\rho_\varepsilon(z)} \quad \forall z \in \mathbb{C}. \quad (3)$$

Proof: By the mean value inequality (42),

$$\begin{aligned} g(z) &\leq \int_{D(z,\varepsilon)} g(w) dm(w) \\ &= \int_{D(z,\varepsilon)} g(w) \frac{1}{\rho(w)} \rho(w) dm(w) \\ &\leq \frac{1}{\rho_\varepsilon(z)} \int_{D(z,\varepsilon)} g(w) \rho(w) dm(w) \\ &= \frac{1}{\pi \varepsilon^2} \frac{1}{\rho_\varepsilon(z)} \|g\|_{L^1(\lambda)}. \end{aligned}$$

For Eq. (3), apply (2) with $g(z) := |f(z)|^2$.

Q.E.D.

Liouville's Theorem

Definition 4. Let $\mathcal{HP}(\mathbb{C})$ denote the space of holomorphic polynomials. Further let,

$$\mathcal{HP}_k = \{p \in \mathcal{HP}(\mathbb{C}) : \deg(p) \leq k\} = \left\{ p(z) = \sum_{n=0}^k a_n z^n : a_n \in \mathbb{C} \right\}.$$

Corollary 5 (Liouville's Theorem). *Suppose there exists $c < \infty$ and $n \in \mathbb{N}_0$ such that*

$$\rho(z) \geq \frac{c}{|z|^{2n} + 1} \text{ for all } z \in \mathbb{C}.$$

Then $\mathcal{HL}^2(\rho) = \mathcal{HP}_k$ for some $k < n$ where $\mathcal{HP}_k := \{0\}$ if $k \leq 0$.

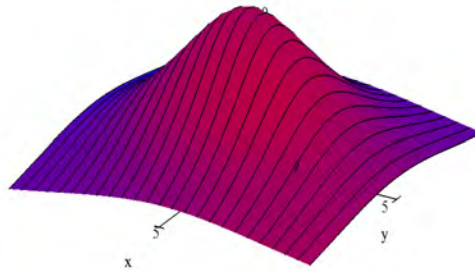


Figure 1: Here, $\rho(z) = 10 / (1 + 0.1 \cdot |z|^2)$ and $\mathcal{HL}^2(\rho) \cong \mathbb{C}$.

Proof: Use the pointwise bounds along with Cauchy estimates (see 44).

Q.E.D.

A Non-Uniform Decay Example

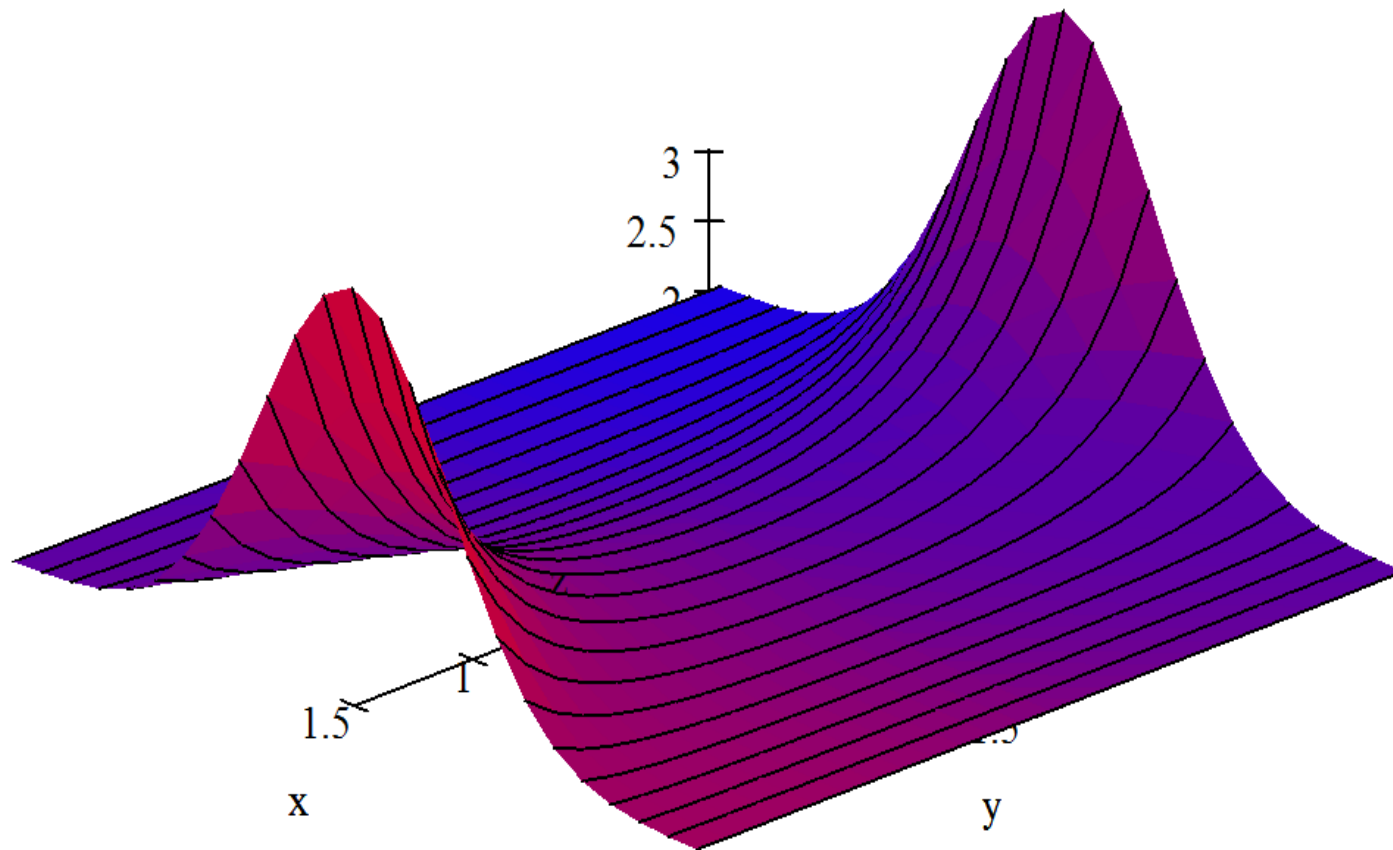


Figure 2: Plot of $d\lambda_c/dm$ for $c = 2$.

- For $c > 0$, let

$$d\lambda_c(z) = \frac{1}{\pi} \exp\left(-\left((1-c)x^2 + (1+c)y^2\right)\right) dx dy.$$

- If $c > 1$, $\mathcal{H}L^2(\lambda_c)$ does **not** contain any polynomials other than 0.
- Nevertheless, $\dim \mathcal{H}L^2(\lambda_c) = \infty$ since

$$\mathcal{H}L^2(\lambda_0) \ni f(z) \rightarrow e^{\frac{c}{2}z^2} f(z) \in \mathcal{H}L^2(\lambda_c)$$

is unitary.

$\mathcal{H}L^2(\lambda)$ is a Hilbert Space

Theorem 6 ($\mathcal{H}L^2(\lambda)$ is Hilberitan). $\mathcal{H}L^2(\lambda)$ is a closed subspace of $L^2(\lambda)$ and hence is a **Hilbert space**.

Proof: For $\varepsilon, r > 0$ let

$$C(r, \varepsilon) := \frac{1}{\pi \varepsilon^2} \cdot \sup_{|w| \leq r} \frac{1}{\rho_\varepsilon(w)} = \frac{1}{\pi \varepsilon^2} \cdot \frac{1}{\min_{|w| \leq r+\varepsilon} \rho_\varepsilon(w)}.$$

Then by the pointwise bounds,

$$\sup_{|w| \leq r} |f(w)|^2 \leq C(r, \varepsilon) \|f\|_{L^2(\lambda)}^2 \text{ for all } f \in \mathcal{H}L^2(\lambda).$$

So if $\{f_n\}_{n=1}^\infty \subset \mathcal{H}L^2(\lambda)$ and $f_n \rightarrow f \in L^2(\lambda)$, we have,

$$\sup_{|w| \leq r} |f_n(w) - f_m(w)|^2 \leq C(r, \varepsilon) \|f_n - f_m\|_{L^2(\lambda)}^2 \rightarrow 0 \text{ as } m, n \rightarrow \infty.$$

So $\{f_n\}_{n=1}^\infty$ is uniformly convergent on compacts and therefore $f \in \mathcal{H}(\mathbb{C})$. **Q.E.D.**

The Reproducing Kernel

Theorem 7. *There exists a unique function, $k(z, w) = k_\lambda(z, w) \in \mathbb{C}$ such that for all $w \in \mathbb{C}$, there exists a unique $k(\cdot, w) \in \mathcal{H}L^2(\lambda)$ such that*

$$f(w) = (f, k(\cdot, w))_{L^2(\lambda)} \quad \forall f \in \mathcal{H}L^2(\lambda). \quad (4)$$

Moreover (see 45)

1. $k(w, z) = (k(\cdot, z), k(\cdot, w))$ and hence $\overline{k(w, z)} = k(z, w)$.
2. $k(z, \bar{w})$ is a holomorphic function of (z, w) .
3. If $\{\varphi_n\}_{n=0}^\infty \subset \mathcal{H}L^2(\lambda)$ is any orthonormal basis, then

$$k(z, w) = \sum_{n=0}^{\infty} \varphi_n(z) \overline{\varphi_n(w)}. \quad (5)$$

The sum is absolutely convergent.

4. For all $w, z \in \mathbb{C}$,

$$\|k(\cdot, z)\|_{L^2(\lambda)}^2 = k(z, z) \leq \frac{1}{\pi \varepsilon^2} \frac{1}{\rho_\varepsilon(z)} \text{ and}$$

$$|k(z, w)| \leq \sqrt{k(z, z) \cdot k(w, w)} \leq \frac{1}{\pi \varepsilon^2} \frac{1}{\sqrt{\rho_\varepsilon(z) \cdot \rho_\varepsilon(w)}}.$$

Optimal Pointwise Bounds

Corollary 8 (Optimal Pointwise Bounds). For all $f \in \mathcal{H}L^2(\lambda)$,

$$|f(w)|^2 \leq k(w, w) \|f\|_{L^2(\lambda)}^2 \text{ for all } w \in \mathbb{C}.$$

These pointwise bounds are optimal.

Proof: By the Cauchy-Schwarz inequality,

$$\begin{aligned} |f(w)|^2 &= \left| (f, k(\cdot, w))_{L^2(\lambda)} \right|^2 \\ &\leq \|k(\cdot, w)\|_{L^2(\lambda)}^2 \|f\|_{L^2(\lambda)}^2 = k(w, w) \|f\|_{L^2(\lambda)}^2. \end{aligned}$$

The function $f(z) := k(z, w)$ saturates this inequality.

Q.E.D.

The Radial Symmetric Case

Theorem 9. Suppose that $\rho(z) = \rho(|z|)$ and $\mathcal{HP} \subset \mathcal{HL}^2(\lambda)$, i.e.

$$a_n^2 := \int_{\mathbb{C}} |z|^{2n} \rho(z) dm(z) < \infty \text{ for all } n \in \mathbb{N}_0.$$

Then

1. $\left\{ \frac{z^n}{a_n} \right\}_{n=0}^{\infty}$ forms an orthonormal basis for $\mathcal{HL}^2(\lambda)$.

2. For any $f \in \mathcal{HL}^2(\lambda)$,

$$f(z) = \sum_{n=0}^{\infty} \frac{f^{(n)}(0)}{n!} z^n$$

converges pointwise and $L^2(\lambda)$.

Proof

If $f \in \mathcal{H}L^2(\lambda)$, then (using Taylor's theorem to evaluate the angular integral)

$$\begin{aligned}(f, z^n) &= \int_0^\infty \left(\int_{-\pi}^\pi f(re^{i\theta}) r^n e^{-in\theta} d\theta \right) \rho(r) r dr \\ &= \int_0^\infty \left(2\pi r^{2n} \frac{f^{(n)}(0)}{n!} \right) \rho(r) r dr = a_n^2 \frac{f^{(n)}(0)}{n!}.\end{aligned}$$

From this it follows that

$$\left\{ \frac{z^n}{a_n} \right\}_{n=0}^\infty \text{ is orthonormal subset of } \mathcal{H}L^2(\lambda).$$

Let $P : \mathcal{H}L^2(\lambda) \rightarrow \mathcal{H}L^2(\lambda)$ be orthogonal projection onto $\overline{\mathcal{HP}}$. Then

$$Pf = \sum \frac{1}{a_n^2} (f, z^n) z^n = \sum_{n=0}^\infty \frac{f^{(n)}(0)}{n!} z^n \quad \forall f \in \mathcal{H}L^2(\lambda)$$

converges in $L^2(\lambda)$ and pointwise to f (by Taylor's theorem) and so $f = Pf \in \overline{\mathcal{HP}}$.

Density of Polynomials

Corollary 10 (Density of Polynomials). When $\rho(z) = \rho(|z|)$, \mathcal{HP} is dense in $\mathcal{HL}^2(\lambda)$.

Proof: See the above proof or see 47 for an alternate proof.

Q.E.D.

Question. Under what conditions on ρ is \mathcal{HP} is dense in $\mathcal{HL}^2(\lambda)$?

Remark. We know (see 53 or 10) \mathcal{HP} is dense in $\mathcal{HL}^2(\lambda)$ if $\rho(z) = \tilde{\rho}(|az + b|)$ for some $a \neq 0$. It is also true if

$$\rho(z) = C \exp\left(-\left(ax^2 + 2bxy + cy^2\right)\right)$$

for some $a, b > 0$ and $c \in \mathbb{R}$ such that $b^2 - ac < 0$.

Radial Symmetric Case Summary

Notation. The **Taylor map** is: $Tf := \alpha \in \mathcal{D}$, where $\alpha_n := f^{(n)}(0)$. Let,

$$a_n^2 := \int_{\mathbb{C}} |z|^{2n} d\lambda(z), \quad \|\alpha\|_{\rho}^2 := \sum_{n=0}^{\infty} |\alpha_n|^2 \left(\frac{a_n}{n!}\right)^2, \text{ and}$$

$$J(\lambda) := \left\{ \alpha = (\alpha_0, \alpha_1, \dots) \in \mathbb{C}^{\mathbb{N}_0} : \|\alpha\|_{\rho}^2 < \infty \right\}.$$

Theorem 11 (Radial Case). *If $\rho(z) = \rho(|z|)$, then $T : \mathcal{H}L^2(\lambda) \rightarrow J(\lambda)$ is **unitary**.*

Moreover, for all $f \in \mathcal{H}(\mathbb{C})$,

$$\int_{\mathbb{C}} |f(z)|^2 \rho(z) dm(z) = \sum_{n=0}^{\infty} \left| f^{(n)}(0) \right|^2 \left(\frac{a_n}{n!}\right)^2 \quad (\text{Isometry Property.})$$

and

$$|f(z)|^2 \leq \|f\|_{L^2(\lambda)}^2 \left(\sum_{n=0}^{\infty} \frac{1}{a_n^2} |z|^{2n} \right). \quad (\text{Optimal Pointwise Bounds.})$$

$$k(z, w) = k_{\lambda}(z, w) = \sum_{n=0}^{\infty} \frac{1}{a_n^2} (z\bar{w})^n \quad (\text{Reproducing Kernel.})$$

Proof (Skip)

The fact that $T : \mathcal{H}L^2(\lambda) \rightarrow J(\lambda)$ is unitary is a translation of the fact that $\left\{ \frac{z^n}{a_n} \right\}_{n=0}^{\infty}$ is an orthonormal basis for $\mathcal{H}L^2(\lambda)$ and the identity,

$$(f, z^n) = a_n^2 \frac{f^{(n)}(0)}{n!}.$$

To see the isometry property is valid for all $f \in \mathcal{H}(\mathbb{C})$, use $T : \mathcal{H}L^2(\lambda) \rightarrow J(\lambda)$ is unitary, Taylor's theorem, and **Fatou's lemma**, to show;

$$\begin{aligned} \int_{\mathbb{C}} |f(z)|^2 \rho(z) dm(z) &= \int_{\mathbb{C}} \liminf_{N \rightarrow \infty} \left| \sum_{n=0}^N \frac{f^{(n)}(0)}{n!} z^n \right|^2 \rho(z) dm(z) \\ &\leq \liminf_{N \rightarrow \infty} \int_{\mathbb{C}} \left| \sum_{n=0}^N \frac{f^{(n)}(0)}{n!} z^n \right|^2 \rho(z) dm(z) \\ &= \liminf_{N \rightarrow \infty} \sum_{n=0}^N a_n^2 \left| \frac{f^{(n)}(0)}{n!} \right|^2 = \sum_{n=0}^{\infty} a_n^2 \left| \frac{f^{(n)}(0)}{n!} \right|^2. \end{aligned}$$

Exponential Examples

Notation. For $\kappa > 0$, let

$$\rho_\kappa(z) := \frac{\kappa}{2\pi} \exp(-|z|^\kappa) \quad \text{and} \quad \Gamma(z) := \int_0^\infty t^z e^{-t} \frac{dt}{t}$$

Theorem 12. If $\rho = \rho_\kappa$, then

$$a_n^2 = \Gamma\left(\frac{2n+2}{\kappa}\right), \quad k(z, w) = \sum_{n=0}^{\infty} \frac{1}{\Gamma\left(\frac{2n+2}{\kappa}\right)} (z\bar{w})^n,$$

and for all $f \in \mathcal{H}(\mathbb{C})$,

$$\int_{\mathbb{C}} |f(z)|^2 \frac{\kappa}{2\pi} \exp(-|z|^\kappa) dm(z) = \sum_{n=0}^{\infty} \left| f^{(n)}(0) \right|^2 \frac{\Gamma\left(\frac{2n+2}{\kappa}\right)}{(n!)^2}$$

and

$$|f(z)|^2 \leq \|f\|_{L^2(\rho_\kappa dm)}^2 \left(\sum_{n=0}^{\infty} \frac{|z|^{2n}}{\Gamma\left(\frac{2n+2}{\kappa}\right)} \right).$$

Example ($\kappa = 1$)

$$k(z, w) = \sum_{n=0}^{\infty} \frac{1}{(2n+1)!} (z\bar{w})^n = \frac{\sinh(\sqrt{z\bar{w}})}{\sqrt{z\bar{w}}}$$

For all $f \in \mathcal{H}(\mathbb{C})$,

$$\frac{1}{2\pi} \int_{\mathbb{C}} |f(z)|^2 \exp(-|z|) dm(z) = \sum_{n=0}^{\infty} \left| f^{(n)}(0) \right|^2 \frac{(2n+1)!}{(n!)^2},$$

and

$$|f(z)|^2 \leq \|f\|_{L^2(\lambda)}^2 \frac{\sinh(|z|)}{|z|} \leq \|f\|_{L^2(\lambda)}^2 \frac{1}{2|z|} e^{|z|}.$$

Example ($\kappa = 2$)

$$d\lambda(z) = \frac{1}{\pi} \exp(-|z|^2) dm(z)$$

$$k(z, w) = \sum_{n=0}^{\infty} \frac{1}{n!} (z\bar{w})^n = e^{z\bar{w}}.$$

For all $f \in \mathcal{H}(\mathbb{C})$,

$$\frac{1}{\pi} \int_{\mathbb{C}} |f(z)|^2 \exp(-|z|^2) dm(z) = \sum_{n=0}^{\infty} \frac{1}{n!} \left| f^{(n)}(0) \right|^2,$$

and

$$|f(z)|^2 \leq \|f\|_{L^2(\lambda)}^2 e^{|z|^2}. \quad \text{(Bargmann's Pointwise Bounds)}$$

References: V.A. Fock (1932) [Fock, 1928], Segal [Segal, 1956, Segal, 1962] and Bargmann [Bargmann, 1961]. (See also Gross and Malliavin [Gross & Malliavin, 1996] for more history.)

Heat Kernel Interpretation for $\kappa = 2$

Fact.

$$\left(e^{t\Delta/4} g \right) (z) = \int_{\mathbb{C}} \frac{1}{\pi t} \exp \left(- |z - w|^2 / t \right) g (w) dm (w).$$

In particular taking $t = 1$ and $z = 0$ implies,

$$\int_{\mathbb{C}} |f (z)|^2 \frac{1}{\pi} \exp \left(- |z|^2 \right) dm (z) = \left(e^{\Delta/4} |f|^2 \right) (0).$$

Recalling that $\Delta = 4\partial\bar{\partial}$ and that $\partial\bar{\partial} |f|^2 = |\partial f|^2$, we have formally,

$$e^{\Delta/4} |f|^2 = e^{\partial\bar{\partial}} |f|^2 = \sum_{n=0}^{\infty} \frac{1}{n!} (\partial\bar{\partial})^n |f|^2 = \sum_{n=0}^{\infty} \frac{1}{n!} |\partial^n f|^2.$$

Combining these last two equations explains why (in this case) that

$$\int_{\mathbb{C}} |f (z)|^2 \frac{1}{\pi} \exp \left(- |z|^2 \right) dm (z) = \sum_{n=0}^{\infty} \frac{1}{n!} \left| f^{(n)} (0) \right|^2.$$

The Segal – Bargmann Transform

Theorem 13 (The Segal – Bargmann isometry). *For all $f \in L^2(\mathbb{R}, d\lambda)$,*

$$\frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} |f(x)|^2 e^{-x^2/2} dx = \frac{1}{\pi} \int_{\mathbb{C}} \left| \left(e^{\frac{1}{2}\partial_x^2} f \right)_a(z) \right| \exp\left(-\frac{1}{4}|z|^2\right) dm(z).$$

Also see 52.

Proof: Let us recall,

$$\partial_x^2 = (\partial + \bar{\partial})^2 = \partial^2 + \bar{\partial}^2 + 2\partial\bar{\partial}.$$

By density of $\mathcal{H}L^2(\mathbb{C}, \lambda)$ in $L^2(\mathbb{R}, \lambda)$, it suffices to assume $f \in \mathcal{H}L^2(\mathbb{C}, \lambda)$. In this case,

$$\begin{aligned} \frac{1}{\sqrt{2\pi}} \int_{\mathbb{R}} |f(x)|^2 e^{-x^2/2} dx &= \left(e^{\frac{1}{2}\partial_x^2} |f|^2 \right) (0) = \left(e^{\frac{1}{2}[\partial^2 + \bar{\partial}^2 + 2\partial\bar{\partial}]} |f|^2 \right) (0) \\ &= \left(e^{\partial\bar{\partial}} e^{\frac{1}{2}\partial^2} e^{\frac{1}{2}\bar{\partial}^2} [f \cdot \bar{f}] \right) (0) = e^{\partial\bar{\partial}} \left(e^{\frac{1}{2}\partial^2} f \cdot e^{\frac{1}{2}\bar{\partial}^2} \bar{f} \right) (0) \\ &= e^{\Delta_{\mathbb{C}}/4} \left(\left| e^{\frac{1}{2}\partial^2} f \right|^2 \right) (0) \\ &= \frac{1}{\pi} \int_{\mathbb{C}} \left| \left(e^{\frac{1}{2}\partial_x^2} f|_{\mathbb{R}} \right)_a(z) \right| \exp\left(-\frac{1}{4}|z|^2\right) dm(z). \end{aligned}$$

Q.E.D.

Generalizations to Lie Groups

- $G =$ complex simply connected Lie group (e.g. $SL(n, \mathbb{C})$)
- $\mathfrak{g} = T_e G$ its Lie algebra (e.g. $sl(n, \mathbb{C})$)
- \mathfrak{g}^* = the dual space of \mathfrak{g}
- $q =$ a non-negative Hermitian form on \mathfrak{g}^* (e.g. $q(A, B) = \text{tr}(B^* A)$)

Fact. There exists $m \leq \dim_{\mathbb{C}}(\mathfrak{g})$ and a linearly independent set, $\{X_l\}_{l=1}^m$, such that

$$q(\alpha, \beta) = \sum_{l=1}^m \alpha(X_l) \overline{\beta(X_l)}$$

for all $\alpha, \beta \in \mathfrak{g}^*$.

Definition 14 (Horizontal subspace). The horizontal subspace associated to q is $H = H(q) := \text{span}(X_l : 1 \leq l \leq m)$ with the inner product: $(X_l, X_k)_H := \delta_{lk}$.

Derivative Spaces

- $q^{\otimes k}$ = the extension of q to $(\mathfrak{g}^{\otimes k})^*$, i.e.

$$q^{\otimes k}(\alpha) = \sum_{l_1, \dots, l_k=1}^m |\alpha(X_{l_1} \otimes \dots \otimes X_{l_k})|^2$$

- $T(\mathfrak{g})$ is the tensor algebra over \mathfrak{g} and $T(\mathfrak{g})' = \prod_{k=0}^{\infty} (\mathfrak{g}^{\otimes k})^*$.

- For each $t > 0$ define

$$q_t = \sum_{k=0}^{\infty} \frac{t^k}{k!} q^{\otimes k}$$

(6)

- $J = \langle \xi \otimes \eta - \eta \otimes \xi - [\xi, \eta] : \xi, \eta \in \mathfrak{g} \rangle \subset T(\mathfrak{g})$
- $J^0 = \{ \alpha \in T(\mathfrak{g})' : \alpha|_J \equiv 0 \}$ – the “Derivative Space.”

-

$$J_t^0 := \{ \alpha \in J^0 : q_t(\alpha) < \infty \}.$$

Two Algebraic Theorems

Definition 15 (*Hörmander's condition*). We say q satisfies *Hörmander's condition* if $\text{Lie}(H(q)) = \mathfrak{g}$.

Theorem 16 (D., Gross, Saloff-Coste). *The following are equivalent:*

1. *Hörmander's condition holds, i.e. $\text{Lie}(H) = \mathfrak{g}$.*
2. $T(\mathfrak{g}) = T(H) + J$.
3. *for any $t > 0$, $q_t|_{J_t^0}$ is an inner product on J_t^0 .*

Theorem 17 (D., Gross, Saloff-Coste). *If \mathfrak{g} is "stratified," then the finite rank tensors in J^0 are dense in J_t^0 .*

Remark. For general \mathfrak{g} there are typically **no** finite rank tensors in J^0 , see [Gross, 1998].

The Heat Kernel

- $\tilde{A}(g) = L_{g*}A$ for all $A \in \mathfrak{g}$ and $g \in G$

- **(Laplacian)** $\Delta = \Delta_q := \sum_{l=1}^m \left[\tilde{X}_l^2 + \left(i\tilde{X}_l \right)^2 \right]$

- **(Heat Kernel)** Let $\rho_t : G \rightarrow (0, \infty)$ satisfy,

$$\left(e^{t\Delta/4} f \right) (e) = \int_G f(g) \rho_t(g) dg.$$

where dg denotes a *right Haar* measure on G .

Fact. The heat kernel, ρ_t , satisfies:

$$\begin{cases} \partial \rho_t(x, \cdot) / \partial t = (1/4) \Delta \rho_t(x, \cdot) \\ \rho_t(x, y) dy \rightarrow \delta_x(dy) \text{ (weakly) as } t \rightarrow 0. \end{cases} \quad (7)$$

$\rho_t \in C^\infty(G, (0, \infty))$ by Hörmander's theorem [Hörmander, 1967].

The Taylor Isomorphism

- Let \mathcal{H} = the holomorphic functions on G
- For $\beta = A_1 \otimes \cdots \otimes A_n \in T(\mathfrak{g})$, let $\tilde{\beta} = \tilde{A}_1 \dots \tilde{A}_n$
- For $f \in \mathcal{H}$ and $x \in G$, let

$$\langle \hat{f}(x), \beta \rangle = (\tilde{\beta}f)(x) \quad \forall \beta \in T(\mathfrak{g}). \quad (8)$$

- $\hat{f}(x) \in J^0$ is the **Taylor coefficient** at x .

- *Taylor map* $(Tf := \hat{f}(e))$,

$$\mathcal{H} \cap L^2(G, \rho_t) \ni f \xrightarrow{T} \hat{f}(e) \in J_t^0. \quad (9)$$

Theorem 18 (D., Gross, Saloff-Coste). *If G is simply connected and q satisfies Hörmander's condition, then the **Taylor map, $T : \mathcal{H}L^2(\rho_t) \rightarrow J_t^0$ is unitary.** Moreover,*

$$\int_G |f(g)|^2 \rho_t(g) dg = \left\| \hat{f}(e) \right\|_t^2 \quad \text{for all } f \in \mathcal{H}(G).$$

The “Classical” Example

- $G = \mathbb{C}^d$ with additive group structure
- $H = \mathfrak{g} = \mathbb{C}^d$, $X_l = e_l$ for $l = 1, 2, \dots, d = m$
- $q(\alpha) = \sum_{l=1}^d |\alpha(e_l)|^2$
- $d(w, z) = |z - w|$
- $\Delta = \sum_{l=1}^d \left(\frac{\partial^2}{\partial x_l^2} + \frac{\partial^2}{\partial y_l^2} \right)$ where $z = x + iy$.

$$\rho_t(z) = \left(\frac{1}{\pi t} \right)^d \exp \left(-|z|^2 / t \right)$$

- $J^0 =$ Symmetric Tensors **= Bosonic Fock Space**
- For $f \in \mathcal{H}$, $\hat{f}(z) \in J^0$ since mixed partial derivatives commute.
- References: V.A. Fock (1932) [Fock, 1928], Segal [Segal, 1956, Segal, 1962] and Bargmann [Bargmann, 1961]. (See also Gross and Malliavin [Gross & Malliavin, 1996] for more history.)
- For proofs, go to 54 and 55.

Some History

The Taylor Isomorphism Theorem 18 was known to hold for *non-degenerate* q in the following cases:

1. $G = K_{\mathbb{C}}$: Driver [Driver, 1995] (inspired by B. Hall [Hall, 1994])
2. G arbitrary: Driver and Gross [Driver & Gross, 1997]
3. $G =$ infinite dimensional complex Hilbert-Schmidt orthogonal group: M. Gordina, [Gordina, 2000b] and [Gordina, 2000a]
4. $G =$ invertible operators in a factor of type II_1 : M. Gordina in [Gordina, 2002]
5. $G =$ path and loop groups of a “stratified” Lie group: M. Cecil, in [Cecil, 2006].
6. $G =$ infinite dimensional Heisenberg like groups, see [Driver & Gordina, 2008b, Driver & Gordina, 2008a, Driver & Gordina, 2008c].
7. For the case presented here see, [Driver *et al.*, 2009b], [Driver *et al.*, 2009c], and [Driver *et al.*, 2009a].

Isometry Proof

- Working analogously to the “ $\kappa = 2$ example” above one sees, **formally**, that

$$\|f\|_{L^2(\rho_t dm)}^2 = e^{t\Delta/4} |f|^2(e) = \sum_{n=0}^{\infty} \frac{1}{n!} \sum_{l_1, \dots, l_k=1}^m |\tilde{X}_{l_1} \dots \tilde{X}_{l_k} f(e)|^2 = \|\hat{f}\|_t^2.$$

- To make this rigorous takes a fair amount of work and requires:
 1. Gaussian heat kernel bounds which involve the “Carnot-Caratheodory” distance on G associated to q (see 56).
 2. Good a-priori pointwise bounds for f and there derivatives.
 3. Careful attention to the fact that finite rank tensors are not dense in $J_t^0(\mathfrak{g})$ in general.
 4. Similarly we must deal with the complication of not knowing a simple to use dense subset of $\mathcal{H}L^2(\rho_t dm)$.

Surjectivity Proof

- The surjectivity proof require the reconstruction of a holomorphic function from its derivatives, $\alpha \in J_t^0(\mathfrak{g})$.

Notation (Rolling Map). Associated to a finite energy path, $g : [0, 1] \rightarrow G$, from e to $z \in G$, let

$$b(s) = b(g, s) := \int_0^s L_{g(t)^{-1}*} \dot{g}(t) dt \in \mathfrak{g}.$$

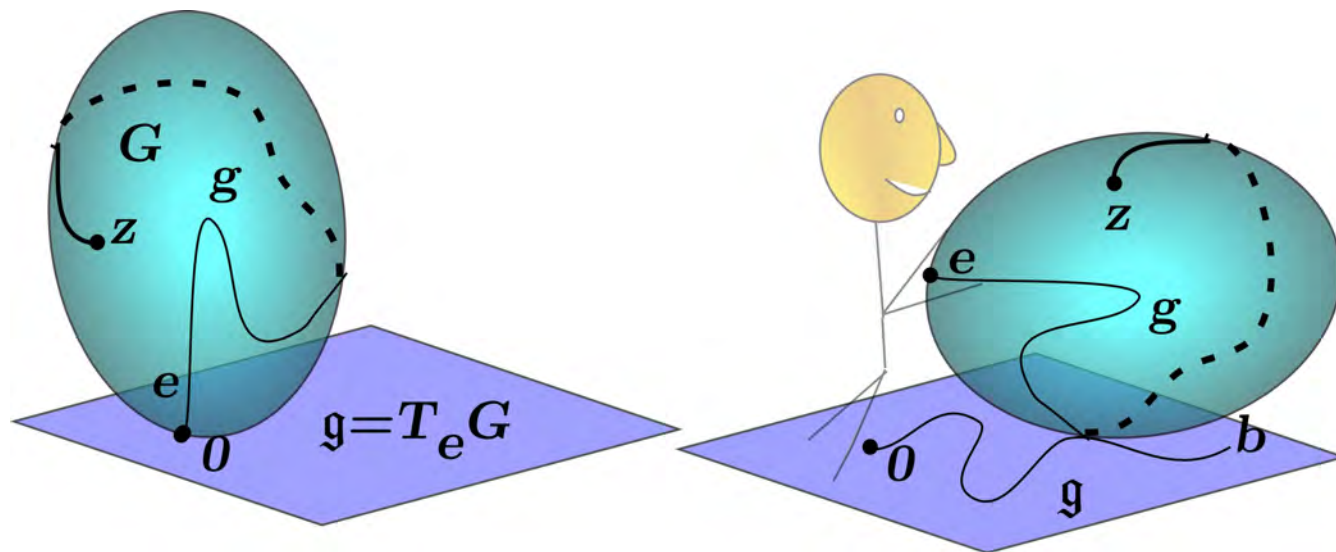


Figure 3: Cartan's rolling map.

Group Taylor Series

Theorem 19 (A **Reconstruction Theorem**). *Suppose;*

1. $g : [0, 1] \rightarrow G$ such that $g(0) = e$ and $g(1) = z$,
2. $f \in \mathcal{H}$ or f is holomorphic near $e \in G$,
3. $\Psi(g) := \sum_{n=0}^{\infty} \Psi_n(g)$ where

$$\Psi_n(g) := \int_{0 \leq s_1 \leq \dots \leq s_n \leq 1} db(s_1) \otimes db(s_2) \otimes \dots \otimes db(s_n).$$

Then

$$f(z) = \langle \hat{f}(e), \Psi(g) \rangle = \sum_{n=0}^{\infty} \langle \hat{f}(e), \Psi_n(g) \rangle \quad (10)$$

and if g is horizontal, i.e. $b(s) \in H$, we have the **pointwise bounds**,

$$|f(z)|^2 \leq \|\hat{f}(e)\|_t^2 e^{d_H^2(e,z)/t} \leq \|\hat{f}(e)\|_t^2 e^{\ell_H^2(g)/t} \quad (11)$$

An Exponential Path Example

Suppose that $b(s) = sA$ for some $A \in \mathfrak{g}$. Then

$$g(s) = e^{sA}$$

and Eq. (10) reduces to the familiar formula,

$$f(e^A) = \left\langle \hat{f}(e), \sum_{n=0}^{\infty} \frac{1}{n!} A^{\otimes n} \right\rangle = \sum_{n=0}^{\infty} \frac{1}{n!} (\tilde{A}^n f)(e).$$

Proof

For $b \in H^1(\mathfrak{g})$ (the **finite energy paths** in \mathfrak{g}) let $g_t(b)$ solve (see Figure 3)

$$\dot{g}_t(b) = L_{g_t(b)*} \dot{b}(t) \text{ with } g_0(b) = e.$$

1. The map $H^1(\mathfrak{g}) \ni b \rightarrow g_1(b) \in G$ is holomorphic.
2. The map $H^1(\mathfrak{g}) \ni b \rightarrow f(g_1(b)) \in \mathbb{C}$ is holomorphic.
3. By **Taylor's Theorem**,

$$f(g_1(b)) = \sum_{n=0}^{\infty} \frac{1}{n!} \partial_b^n (f \circ g_1)(0).$$

4. A direct but involved computation shows,

$$\frac{1}{n!} \partial_b^n (f \circ g_1)(0) = \langle D^n f(e), \Psi_n(g) \rangle = \langle \hat{f}(e), \Psi_n(g) \rangle.$$

5. The pointwise bounds in (11) follow from (10) and the Cauchy-Schwarz inequality.

Horizontal Reconstruction Theorem

Theorem 20 (Horizontal Reconstruction). Given $\alpha \in J_t^0$, there exists $f \in \mathcal{H}$ such that $\hat{f}(e) = \alpha$.

Proof Ideas

- We must define f by

$$f(g(1)) := \langle \alpha, \Psi(g) \rangle \quad (12)$$

for all paths, g , such that $g(0) = e$.

- However, in the degenerate case, we only know *a priori* that $\langle \alpha, \Psi(g) \rangle$ is well defined when g is *horizontal*.
- How do we show $g \rightarrow \langle \alpha, \Psi(g) \rangle$ only depends on $g(1)$?
- **Answer:** we first construct local version of f and then use an analytic continuation argument to patch them together.

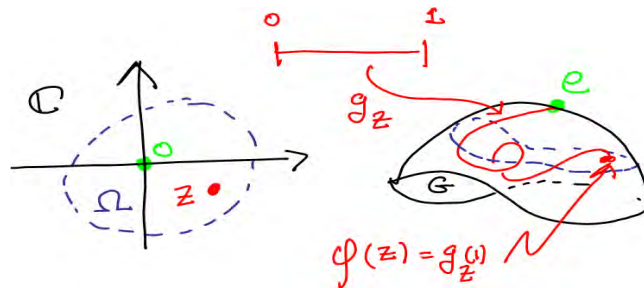
Local Reconstruction Theorem

Theorem 21 (Local Reconstruction). *There exists open neighborhoods, $0 \in \Omega \subset \mathbb{C}^d$ and $e \in U \subset G$ such that:*

1. *for $z \in \Omega$ there exists a horizontal paths, $g_z(t) \in G$, depending holomorphically on z , such that*
2. *if $\varphi(z) := g_z(1)$, then $\varphi : \Omega \rightarrow U$ is biholomorphic.*
3. *The function $f : U \rightarrow \mathbb{C}$ defined by*

$$f(\varphi(z)) := \langle \alpha, \Psi(g_z) \rangle$$

is holomorphic and $\hat{f}(e) = \alpha$.



Example: Complex Heisenberg Group

$G = \mathbb{C}^3$ with group law;

$$\begin{aligned} & (z_1, z_2, z_3) \cdot (z'_1, z'_2, z'_3) \\ &= \left(z_1 + z'_1, z_2 + z'_2, z_3 + z'_3 + \frac{1}{2}(z_1 z'_2 - z_2 z'_1) \right). \end{aligned}$$

- $\mathfrak{g} = \mathbb{C}^3$, $H = \mathbb{C}^2 \times \{0\}$, $X_l = e_l$ for $l = 1, 2$.

- $q(\alpha) = \sum_{l=1}^2 |\alpha(e_l)|^2$

- $\Delta = \tilde{X}_1^2 + \tilde{X}_2^2 + \tilde{Y}_1^2 + \tilde{Y}_2^2$ where $Y_l = iX_l$.

-

$$\Delta_H = \Delta_{z_1} + \Delta_{z_2} + \frac{|z_1|^2 + |z_2|^2}{4} \Delta_{z_3} + L \frac{\partial}{\partial x_3} + S \frac{\partial}{\partial y_3}$$

- L and S are angular momentum ops. on $\mathbb{C}^2 \times \{0\}$.

Heat Kernel and Horizontal Paths

Theorem 22 (Hypoelliptic Heat Kernel). *The heat kernel for the complex Heisenberg group setup above is given by,*

$$\begin{aligned} \rho_t(z) = & \left(\frac{1}{2\pi} \right)^4 \int_{\mathbb{C}} \frac{|w|^2}{\sinh^2(|w|t/4)} \\ & \times \exp \left(-\frac{1}{4} |w| \coth(|w|t/4) \left(|z_1|^2 + |z_2|^2 \right) \right) \\ & \times e^{i \operatorname{Re}(w\bar{z}_3)} dm(w). \end{aligned}$$

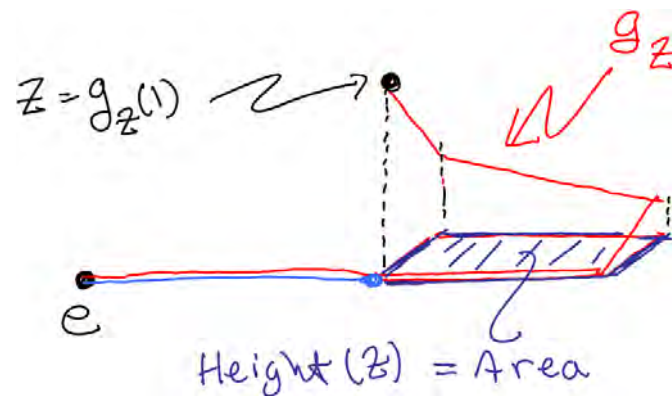


Figure 4: The path g_z for the Heisenberg group. The rectangular region is long and skinny.

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The Taylor map on homogenous spaces

- Let $K \subset G$ be a connected, closed, complex subgroup of G , $\mathfrak{k} = \text{Lie}(K)$.
- $M = K \backslash G$ be the space of right K cosets,
- $\pi : G \rightarrow M$ be the associated quotient map,

Notation. The formula,

$$\dot{A}(m) := \left. \frac{d}{dt} \right|_0 (me^{tA}) \text{ for all } m \in M \text{ and } A \in \mathfrak{g} \quad (13)$$

defines a linear map, $\mathfrak{g} \ni A \rightarrow \dot{A} \in \text{Vect}(M)$, where $\text{Vect}(M)$ denotes the linear space of smooth vector fields on M . We may also define a sub-Laplacian Δ_M on M given by

$$\Delta_M = \sum_{j=1}^m (\dot{X}_j^2 + \dot{Y}_j^2), \text{ where } Y_j := iX_j. \quad (14)$$

Definition 23. Let $\lambda_t(dm)$ be **heat kernel** measure on M given by

$$\lambda_t(dm) = (\pi_* \rho_t)(dm). \quad (15)$$

Theorem 24. *The family of probability measures $\{\lambda_t : t \in (0, \infty)\}$, is the unique family of probability measures on M such that, for all $\phi \in C_c^\infty(M)$, the function $t \rightarrow \lambda_t(\phi) := \int_M \phi d\lambda_t$ is continuously differentiable and satisfies*

$$\frac{d}{dt}\lambda_t(\phi) = \frac{1}{4}\lambda_t(\Delta_M\phi) \quad \text{and} \quad \lim_{t \downarrow 0} \lambda_t(\phi) = \phi(o). \quad (16)$$

Definition 25 (G – space Taylor map). For $u \in \mathcal{H}(M)$, define $\hat{u} \in T'$ by; $\langle \hat{u}, 1 \rangle = u(Ke)$ and for all $n \in \mathbb{N}$,

$$\langle \hat{u}, \xi_1 \otimes \cdots \otimes \xi_n \rangle = (\dot{\xi}_1 \cdots \dot{\xi}_n u)(Ke) = \tilde{\xi}_1 \cdots \tilde{\xi}_n [u \circ \pi](e) \quad \text{for all } \xi_j \in \mathfrak{g}. \quad (17)$$

The map $\mathcal{H}(M) \ni u \rightarrow \hat{u} \in T'$ is called the Taylor map on M .

Theorem 26 (The quotient Taylor map). *For all $t > 0$, the Taylor map*

$$\mathcal{H}(M) \supset \mathcal{HL}^2(M, \lambda_t) \ni u \rightarrow \hat{u} \in (J + \mathfrak{k}T)_t^0 \subset T'$$

is a unitary map, where

$$(J + \mathfrak{k}T)_t^0 = \{\alpha \in T' : \langle \alpha, J + \mathfrak{k}T \rangle = \{0\}\}.$$

Two examples

The Grushin complex 2-space

Notation (Complex Heisenberg group). Let $H_3^{\mathbb{C}} = \mathbb{C}^3$ with the group law

$$(z_1, z_2, z_3) \cdot (z'_1, z'_2, z'_3) = (z_1 + z'_1, z_2 + z'_2, z_3 + z'_3 + (1/2)(z_1 z'_2 - z_2 z'_1)).$$

We take

- Let $q(\alpha) := |\alpha(e_1)|^2 + |\alpha(e_2)|^2$ for all $\alpha \in (\mathbb{C}^3)^*$.
- $K = \mathbb{C} = \{(z_1, z_2, z_3) : z_2 = z_3 = 0\} \subset H_3^{\mathbb{C}}$
- $M = K \backslash H_3^{\mathbb{C}} \cong \mathbb{C}^2 \cong \mathbb{R}^4$; $\xi = (w, z) = (u + iv, x + iy) \in M$,
- It turns out that

$$\Delta_M = (\partial/\partial u)^2 + (\partial/\partial v)^2 + (u^2 + v^2)((\partial/\partial x)^2 + (\partial/\partial y)^2).$$

- The heat kerne density, $\lambda_t(\xi)$ satisfies

$$\frac{c_1}{V(\sqrt{t})} \exp\left(-C_1 \frac{\delta(\xi)^2}{t}\right) \leq \lambda_t(\xi) \leq \frac{C_2}{V(\sqrt{t})} \exp\left(-c_2 \frac{\delta(\xi)^2}{t}\right) \quad (18)$$

- $\delta(\xi)$ is the subelliptic distance between the origin and $\xi = (w, z)$,

$$c(|u| + |v| + \sqrt{|x|} + \sqrt{|y|}) \leq \delta(\xi) \leq C(|u| + |v| + \sqrt{|x|} + \sqrt{|y|}).$$

- $V(r) \simeq r^6, \quad r > 0$

-

$$\Omega(m, n) = m! \cdot \binom{|n|_1}{m_1} \cdot \binom{|n|_2 - |m|_1}{m_2} \cdot \binom{|n|_3 - |m|_2}{m_3} \cdots \binom{|n|_k - |m|_{k-1}}{m_k}$$

Corollary 27. Suppose that f is a holomorphic function on \mathbb{C}^2 , Then

$$\|f\|_t^2 = \|\hat{f}\|_t^2 = \sum_{N=0}^{\infty} \frac{t^N}{N!} \sum_{(m,n) \in I(N)} \Omega^2(m, n) \left| \left(\partial_w^{|n|-|m|} \partial_z^{|m|} f \right) (0, 0) \right|^2. \quad (19)$$

Example. When $f(w, z) = g(z) = z^3$, the only non-zero derivative at $(0, 0)$ is $(\partial^3 f / \partial^3 z)(0, 0) = 6$. So according to Eq. (??),

$$\int_M (x^2 + y^2)^3 \lambda_t(d\xi) = \|\hat{f}\|_t^2 = \frac{61}{20} t^6. \quad (20)$$

Example. For $f(w, z) = wz^3$, we have

$$\left\| \int_M |w|^2 |z|^6 \lambda_t(d\xi) \right\| = \|\hat{f}\|_t^2 = \frac{277}{28} t^7. \quad (21)$$

A one-dimensional complex G - space

Notation. We let $G := \mathbb{C}^\times \ltimes \mathbb{C}$ with group law,

$$(a, b) (a', b') = (aa', ab' + b)$$

Take:

- $q(\alpha) = |\langle \alpha, e_1 \rangle|^2 + |\langle \alpha, e_2 \rangle|^2$ (this is positive!),
- $K = \mathbb{C}^\times \times \{0\}$, $M := K \backslash G \cong \mathbb{C}$.
- $\Delta_M = [1 + x^2 + y^2] (\partial_x^2 + \partial_y^2)$
- The associated heat kernel satisfies,

$$\frac{C_\epsilon}{\ln(\cosh^2 \sqrt{t})} e^{-(1+\epsilon)(\sinh^{-1} |\xi|)^2/t} \leq \lambda_t(\xi) \leq \frac{C_\epsilon}{\ln(\cosh^2 \sqrt{t})} e^{-(1-\epsilon)(\sinh^{-1} |\xi|)^2/t}.$$

(22)

Corollary 28. *Suppose that f is a holomorphic function on \mathbb{C} , then*

$$\int_{\mathbb{C}} |f(\xi)|^2 \lambda_t(\xi) d\xi = \|f\|_t^2 = \sum_{m=0}^{\infty} c_m(t) \left| f^{(m)}(0) \right|^2 \quad (23)$$

where

$$\begin{aligned} c_0(t) &\equiv 1 \\ c_1(t) &= e^t - 1, \\ c_2(t) &= \frac{1}{12}e^{4t} - \frac{1}{3}e^t + \frac{1}{4} \\ c_m(t) &\gtrsim \frac{1}{\sqrt{2\pi t}} \frac{e^{m^2 t}}{m^{2m+1}} \end{aligned}$$